

Solution Structures of Hauser Base ^{*i*}Pr₂NMgCl and *Turbo*-Hauser Base ^{*i*}Pr₂NMgCl·LiCl in THF and the Influence of LiCl on the Schlenk-Equilibrium

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Supporting Information

ABSTRACT: Grignard reagents that are at the simplest level described as "RMgX" (where R is an organic substituent and X a halide) are one of the most widely utilized classes of synthetic reagents. Lately, especially Grignard reagents with amido ligands of the type R_1R_2NMgX , so-called Hauser bases, and their *Turbo* analogue R_1R_2NMgX ·LiCl play an outranging role in modern synthetic chemistry. However, because of their complex solution behavior, where Schlenk-type equilibria are involved, very little is known about their structure in solution. Especially the impact of LiCl on the Schlenk-equilibrium was still obscured by complexity



DOSY-ECC-Molecular Weight-Determination of Hauser Base [/]Pr₂NMgCl and *Turbo*-Hauser Base [/]Pr₂NMgCl·LiCl in THF-d_o solution

and limited analytical access. Herein, we present unprecedented insights into the solution structure of the Hauser base ${}^{i}Pr_{2}NMgCl \cdot LiCl 2$ at various temperatures in THF- d_{8} solution by employing a newly elaborated diffusion ordered spectroscopy (DOSY) NMR method hand-in-hand with theoretical calculations.

INTRODUCTION

Metalation reactions are one of the most important large-scale synthetic methodologies in chemical industry. Prominent reagents for proton abstraction from organic substrates are organolithium reagents like *n*-butyllithium closely followed by sterically demanding amides like diisopropylamide, LDA; 1,1,1,3,3,3-hexamethyldisilazide, LiHMDS; and 2,2,6,6-tetramethyl-piperidide, LiTMP.¹ However, the drawback of these strong organic bases is that they often cause competing side reactions, require relatively low temperatures (e.g., -78 to -90 °C), and do not tolerate certain synthetically important functional groups like esters, carbonyls, nitriles, sulfoxides, and halides. This is why in the past decades organomagnesium reagents have found increasing popularity in synthetic chemistry. As compared to organolithium reagents, the magnesium compounds have more covalent character and therefore less reactive metal-ligand bonds. This is why they display a better functional group tolerance and a much greater chemoselectivity. In the late 1940s, Hauser and co-workers succeeded in the development of the amido Grignard reagents R₂NMgCl, formally known as Hauser bases, by replacing the alkyl ligand of a Grignard reagent with a secondary amide.² The breakthrough in the synthetic application of Hauser bases ensued in the 1980s and 1990s, where Eaton and co-workers introduced ⁱPr₂N- and TMPMgBr, which were shown to orthomagnesiate carboxamides.³ Later, Kondo, Sakamoto, and coworkers reported the utility of Pr_2NMgX (X = Cl, Br) as selective deprotonation reagents for heterocyclic thiophene (exclusively at the 2- position) and phenylsulphonyl-substituted indoles.⁴ In 2006, Knochel and co-workers developed the mixed lithium and magnesium amide bases R1R2NMgCl·LiCl. These so-called Turbo-Hauser bases display a better solubility and a higher reactivity for the magnesiation of various functionalized (hetero)aromatic compounds.⁵ Although there is a great deal of information on the utility of these reagents, very little is known regarding the nature of (Turbo-)Hauser bases in solution. One reason for that paucity is the complex and interlocked equilibrium framework in solution that Hauser bases show. It was proposed that it could be similar to the Schlenk-equilibrium of Grignard reagents in ether solution, where more than one magnesium-containing species exists.⁶ A rearrangement of the organic ligand takes place and ends up in an equilibrium with the diorganomagnesium and the magnesiumdihalide (eq 1 in Scheme 1).⁷ Later, molecular association studies also revealed oligomeric Grignard structures in diethyl ether. This fact further complicated the simple Schlenk-equilibrium (eq 1) as oligomeric species had to be included in an advanced pattern (eq 2 in Scheme 1).⁸ There are a few known solid-state structures of Grignard complexes.⁹ It was found that the Mg atoms are predominantly tetrahedrally coordinated in the solid state and dimeric species are bridged through halide atoms. Especially in the late 1960s, a lot of solution structure investigations have been done mainly by ebullioscopic,^{8,10} calorimetric,¹¹ and NMR-¹² measurements. The position of the Schlenk-equilibrium is considered to be

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dependent on the concentration, the nature of the solvent, the steric demand of the organic substituent, the halide involved, and the temperature.¹³ In THF, all alkyl and arylmagnesium halides (Cl, Br, I) are found to be monomeric over a wide concentration range.^{8,11} In diethyl ether, alkyl and aryl Grignard reagents are mostly monomeric at low concentrations (less than 0.5 M) and mostly dimeric at higher concentrations (0.5-1.0 M).^{10a'} However, to our own surprise, there are no investigations about the aggregation of Hauser bases in solution. Only one study on the solution structure of Turbo-Hauser base 2 and its bulky analogue TMPMgCl LiCl exists.¹⁴ Unfortunately, the exact nature of the solution species could not be identified because the employed diffusion coefficientformula weight (D-FW) analysis¹⁵ that is based on an internal calibration curve (ICC, pioneered by Williard et al.¹⁶) lacks the required accuracy. Many internal standards, which might interfere with the reactive metal complexes, are required. Recently, we developed a novel DOSY-NMR method for accurate molecular weight (MW) determination that rests on external calibration curves (ECC) with normalized diffusion coefficients. One internal reference is sufficient, so the addition of multiple internal references is not required anymore. This method offers independent of diversities in temperature or viscosity an easy adjustment on the analyte geometry and a robust access to determine accurate MWs in solution with an error of better than $\pm 9\%$ (for more information, see the Supporting Information).¹⁷ Using this method, we were already able to characterize the complex oligomeric mixture of LDA in toluene solution.¹⁸ Hence, we embarked to apply our newly developed and improved NMR techniques to identify reactive organometallic species in solution.¹⁹ In cases where an unequivocal assignment is impossible, given the existence of structural isomers, we have complemented the experimental findings with electronic structure calculations. The latter also provide further insight into the aggregation patterns in solution. Understanding how (Turbo-)Hauser bases self-organize and interact among each other, especially in solution where they operate, is a premium in structure-reactivity elucidation.

RESULTS

Solid-State Structure of ^{*i*}**Pr₂NMgCl 1.** To shed light on the solution structure of 1, we first synthesized the prominent Hauser base ^{*i*}Pr₂NMgCl 1 and crystallized it from a THF/ toluene 1:1 mixture at -6 °C (see Experimental Section).²⁰ The crystal structure of 1 is shown in Figure 1. TMP²¹ and HMDS²² Hauser bases alike all Grignard dimers^{9a,b,e,23} show the halides in the bridging position in the solid state (like D2 in category B from Chart 1). In contrast, 1 dimerizes featuring the amido ligands in the bridging position (D1 in category A from Chart 1). Searching the Cambridge Crystallography Database for Hauser bases reveals that there are only three other dimeric



Figure 1. Solid-state structure of $[1 \cdot \text{THF}]_2$ (D1) with hydrogen atoms and disorder omitted for clarity (selected bond lengths and angles are displayed in Figure S1).

Chart 1. Most Plausible Aggregation Modes of ${}^{i}Pr_{2}NMgCl 1$ in THF- d_{8} Solution²⁶



amido bridged Hauser bases in the literature.²⁴ All have one feature in common: less bulky amido ligands like $Et_2N^{-,22a}$ $Ph_3P=N^{-,25}$ and ${}^{i}Pr_2N^{-14}$ in the bridging position. It can be concluded that in the solid state the switch from the halide to the amido bridge seems advantageous.⁶ However, solid-state structures may not necessarily be maintained in solution.

Solution Structure of ⁱPr₂NMgCl 1. In synthesis, organometallic compounds are predominantly used in solution. Therefore, we were interested in the solution structure of 1. The most plausible aggregation modes of a Hauser base in THF solution are filed in Chart 1. A dissolved crystal of 1 in THF can either retain its dimeric status D1 or isomerize to the chloride bridged dimer D2. On the other hand, these dimers can dissociate to the monomer M1 or rearrange according to the Schlenk-equilibrium to the diamidomagnesium M2 and MgCl₂. When an excess of MgCl₂ is present, MgCl₂ salt cocoordinated species like $MI(\mu Cl)_3MgCl_2$, $M2 \cdot MgCl_2$, $MgCl_2 \cdot D1(D2) \cdot MgCl_2$, or $M1 \cdot (MgCl_2)_2 \cdot M1$ might also be present in solution. All of the mentioned species can be distinguished either by their MW or additionally by the chemical environment of the isopropyl groups that is reflected in the chemical shift δ . As compared to terminal amido ligands (category B, Chart 1), the bridging ones (category A, Chart 1) At room temperature, however, the ¹H NMR spectrum of 1 (Figure 2) shows one broad signal set corresponding to a single



Figure 2. Superposition of ¹H NMR spectra of crystalline $[1 \cdot \text{THF}]_2$ (0.10 M, $-\text{CH}_3$ region) redissolved in THF- d_8 at various temperatures. Assignment of signals: see Scheme 2.

species a1 at high field (2.94/1.01 ppm for α -CH/CH₃, Chart 1, category B). The ¹H DOSY-ECC-MW-determination agrees best with the heteroleptic monomer M1 ($MW_{calc} = 304 \text{ g/mol}$, $MW_{det} = 310 \text{ g/mol}, MW_{err} = -2\%)$ ²⁷ while dimeric D2, M2· MgCl₂, and bigger aggregates can be excluded (MW_{err} \geq 33%). Below 0 °C, two additional species **b1** (3.10/1.02 ppm, category B), c1 (3.24/1.01 ppm, category B) at high field and d1 at low field (3.43/1.29 ppm, category A) grow in. From NMR studies on alkyl Grignard reagents it is known that homoleptic dialkylmagnesium monomers analogue M2 resonate at lower field than heteroleptic monomers analogue M1.²⁸ In fact, MW-determination for species **b1** agrees perfectly with the homoleptic diamidomagnesium M2 ($MW_{calc} = 369 \text{ g/mol}$, $MW_{det} = 356$, $MW_{err} = 4\%$). The MW of c1 matches to those of dimeric **D2** and $M2 \cdot MgCl_2$ (MW_{calc} = 464 g/mol, MW_{det} = 450 g/mol, $MW_{err} = 3\%$) that have the same MW. Both have a comparable chemical and magnetic environment and cannot be distinguished by their MWs. Both species could be present in solution, but it is most likely that the equilibrium is significantly shifted to the side of dimer D2 because it displays less steric hindrance as compared to M2·MgCl2. To investigate the structure of the dimer in solution, electronic structure calculations were carried out on the possible isomers. The latter were performed with the B3LYP-D3 method,²⁹ including solvent corrections. Further information can be found in the Computational Details. Free energy differences confirm M2· MgCl₂ to be disfavored relative to the D2 species by 53.1 kJ/ mol (Table S29). The most stable structure found corresponds in fact to a cis-isomer of D2, with both bases orientated to the same side of the Mg₂Cl₂ ring. This arrangement optimizes dispersion interactions between both the propyl moieties and the THF rings on each side. The trans configuration is slightly higher in energy by 7.2 kJ/mol. However, this marginal difference is not to be taken for granted because weak interactions with the solvent (which in our computations is only included as a dielectric continuum) could easily counterbalance this effect. The optimized structure of M2·MgCl₂ shows a large N-Mg-N angle of about 146° (Figure S35), illustrating the steric strain of both diisopropylamido groups

when coordinated to the same metal center adding to the energetic disfavor. Species **d1**, with the highly low field shifted signal, is in good agreement with the amido bridged dimer **D1** ($MW_{calc} = 464 \text{ g/mol}, MW_{det} = 435 \text{ g/mol}, MW_{err} = 6\%$) that is similar to the crystal structure of $[1 \cdot \text{THF}]_2$ (Figure 1). Lowering the temperature dramatically influences the position of the Schlenk-equilibrium of Hauser base **1** (Scheme 2). While



^{*a*}The MWs were derived from ¹H DOSY-ECC-MW-determinations.²⁶ The accuracy of this method is in the range of $MW_{err} < \pm 9\%$.^{17,35}

at room temperature, the monomer M1/M2 ratio is 4:1, this ratio switches completely at -75 °C to the ratio of 1:4 (Table 1, 0.100 M). The huge population of homoleptic monomer M2

Table 1.	Schlenk-Equi	ilibrium "	'Constant"	$K_{\rm s}$ for the	
Reaction	$(^{i}\mathrm{Pr}_{2}\mathrm{N})_{2}\mathrm{Mg}$	(M2) + N	$MgCl_2 \rightleftharpoons 2$	^{<i>i</i>} Pr ₂ NMgCl	(M1)

$T/^{\circ}C$	$K_{\rm s}$ (0.015 M)	$K_{\rm s}$ (0.100 M)
25	25.00	16.00
0	8.16	4.43
-15	2.97	1.49
-25	1.00	0.68
-50	0.18	0.10
-75	0.12	0.05

"Usually the Schlenk-equilibrium constant K_s is derived from ¹H integrals of α -CH protons with $K_s = [{}^{i}Pr_2NMgCl]^2/[({}^{i}Pr_2N)_2Mg]^2$ with the approximation: $[({}^{i}Pr_2N)_2Mg] \approx [MgCl_2]$. However, our results show that the concentration of $({}^{i}Pr_2N)_2Mg$ is not equal to MgCl₂ because the latter is involved in further reactions with Hauser base 1. This is also reflected in our determined K_s values that show therefore concentration dependence.

at low temperatures is in good agreement with the work of Smith and Becker who showed that the formation of RMgCl from R_2Mg and MgCl₂ is endothermic in THF solution.¹¹ The same seems valid for Hauser base 1. The equilibrium constants of the Schlenk-equilibrium summarized in Table 1 display the equilibrium to move to homoleptic $M2 + MgCl_2$ with increasing concentration and decreasing temperature. The formation of M2 is accompanied by the release of free MgCl₂ to the solution, most probably as a monomer with up to four THF molecules.³⁰ Therefore, it might not be surprising when

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free magnesium dihalide coordinates to some complexes at low temperatures. At -50 °C, dimers D1 (d1) and D2 (c1) dissipate at the expense of two new species e1 (3.39/1.50 ppm), category A) and f1 (3.21/1.04 ppm, category B) that display an even more pronounced low field shift. This remarkable shift can be attributed to an additional coordination of MgCl₂ to M1, D1, or D2. When MgCl₂ coordinates to monomer M1, then a structure like $M1(\mu Cl)_3MgCl_2$ (Chart 1, category B) would be feasible. This coordination mode was suggested for methyl magnesium chloride by Sakamoto and Imamoto et al.³¹ With the help of coldspray ionization mass spectrometry (CSI-MS), they proposed that the μ -Cl₃ bridged Grignard reagent was coordinated by four to six THF molecules, whereas the species with five THF molecules was the major component.³¹ Additional support is provided by several crystal structures of cationic $[(THF)_3Mg(\mu Cl)_3Mg(THF)_3]^+$ where magnesium chloride is also coordinated in that μ -Cl₃ fashion.³¹ Our MWinvestigation shows similar results: At -70 °C the ECC-MWdetermination gives for signal f1 a MW of $MW_{det} = 512 \text{ g/mol}$ that fits $M1(\mu Cl)_3MgCl_2$ with four THF molecules (MW_{calc} = 544 g/mol, $MW_{det} = 512$ g/mol, $MW_{err} = 6\%$). At -80 °C the MW increases significantly to $MW_{det} = 616 \text{ g/mol}$ that matches perfectly the THF 5-fold solvated $M1(\mu Cl)_3MgCl_2$ (MW_{calc} = 616 g/mol, $MW_{det} = 616$ g/mol, $MW_{err} = 0\%$), indicating that a higher solvation is favored at lower temperatures.³² In the literature there are several crystal structures known of the type $M1 \cdot (MgCl_2)_2 \cdot M1$ (see Chart 1, category B), where two monomers M1 are bridged by two magnesium dichlorides in an open cubic aggregation mode.^{31,33} However, it seems that the dissociation into smaller parts is favored over that open cubic arrangement in solution of 1 ($MW_{calc} = 943$ g/mol, $MW_{det} = 616 \text{ g/mol}, MW_{err} = 35\%$), which was already proposed by D. Seyferth in 2009.¹³ Additionally to f1, species el appears at much lower field (category A). The ECC-MWdetermination shows a temperature-dependent MW distribution (MW_{det} = 409 g/mol at -50 °C, 446 g/mol at -70 °C, and 578 g/mol at -80 °C). The addition of MgCl₂ to dimers D1 or D2 would produce aggregates like MgCl₂·D1·MgCl₂ or MgCl₂·D2·MgCl₂ (Chart 1, category A). However, several other MgCl₂ coordinated, amido bridged species would also be conceivable. Unfortunately, in the literature there are no crystal structures of MgCl₂ coordinated, ligand bridged Hauser bases or Grignard reagents that could provide a hint to the most plausible aggregates. The high MW of 578 g/mol at -80 °C gives much room for interpretation. This is why we can only speculate how the composition of aggregate e1 could look like. Noticeable is that in contrast to all other species, the $-CH_3$ signal of e1 shifts to lower field with lower temperature (Figure 2). This behavior could be a result of a successive addition of MgCl₂ to D1 or D2. At -80 °C the shift to lower field stops,³² and the MW of e1 matches dimeric MgCl₂·D1(D2)·MgCl₂ $(MW_{calc} = 655 \text{ g/mol}, MW_{det,corr} = 703 \text{ g/mol}, MW_{err} = -7\%,$ after molar density correction (see Supporting Information section XXXIII)).³⁴ Although we still cannot be sure about this MW-agreement, it seems that at low temperatures free, monomeric MgCl₂ is disadvantageous in solution. Instead, it coordinates to monomeric and/or dimeric RMgCl molecules as discussed for LiCl in the following section.

Solution Structure of ^{*i*}**Pr₂NMgCl·LiCl 2.** The impact of LiCl on the solution structure of Grignard reagents and Hauser bases is still vigorously discussed. Knochel et al. suggest that LiCl deaggregates RMgX oligomers^{Sd} and forms a more reactive bimetallic monomer RMgCl·LiCl that is supposed to

furnish magnesiate character to the Grignard reagent in the sense of a solvent separated ion pair (SSIP) $[Li(THF)_4]^+[RMg-(THF)Cl_2]^{-.36}$ Crystallographic evidence of the *Turbo*-Hauser bases TMPMgCl·LiCl²¹ (**M1·LiCl**, Chart 2 with B = TMP⁻)

Chart 2. Most Plausible Aggregation Modes of ${}^{i}Pr_{2}NMgCl$ · LiCl 2 in THF- d_{8} Solution²⁶



and $[2 \cdot \text{THF}]_2^{14}$ (LiCl·D1·LiCl, Chart 2) supports the contact ion pair (CIP) coordination mode in the solid state. Yet still it was not clear whether the mixed metal structure really is maintained in solution or just a transient species. García-Álvarez and Mulvey et al. analyzed¹⁴ crystals of [2.THF]₂ in THF- d_8 solution at -50 °C by employing the ICC D-FW analysis that was pioneered by Li and Williard et al.¹⁶ Because of the lack of appropriate references, the accuracy of the method was not sufficient, and they concluded that they were not able to "clearly establish the exact nature of the solution species".¹⁴ However, the first key conclusion was that the molecular structure of crystalline $[2 \cdot THF]_2$ (LiCl·D1·LiCl) was not retained in THF- d_8 solution, and the second was that a SSIP situation like it was proposed by Knochel et al., described by negative charged magnesium ate complexes (like M0), and free $[Li(THF)_4]^+$ (Li2) seemed most probable (Chart 2, category B).¹⁴ In the following section, we will shed light on the complex solution structure of 2 and prove that both key conclusions have to be revised. Furthermore, we will try to clarify the influence of LiCl on the Schlenk-equilibrium of 1 more generally by applying our accurate DOSY-ECC-MWdetermination method¹⁷ and electronic structure calculations.

Primarily, it seems advisible to a priori rationalize which species are feasible to be present in the solution of ${}^{i}Pr_2NMgCl$ · LiCl 2. A dissolved crystal of $[2 \cdot THF]_2$ in THF can either retain the coordination LiCl·D1·LiCl or isomerize to LiCl· D2*·LiCl (Chart 2, category A). In the first aggregation mode, the isopropyl groups are neighbored by compact chloride ligands, enabling a free rotation of the isopropyl groups, whereas in the latter they are neighbored to a sterically demanding THF molecule that would result in a steric

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repulsion. This is why we suggest that species LiCl·D1·LiCl should be highly favored over LiCl·D2*·LiCl.

This again is supported by our calculations (Table S31). The free energy difference between LiCl·D1·LiCl and LiCl·D2*· LiCl is 55–56 kJ/mol (the former being more stable) in the temperature range from -90 to +25 °C. Inspection of the optimized structures confirms the steric hindrance (Figure \$38). In the LiCl·D1·LiCl structure the Mg²⁺ cations are aligned with the Li⁺, resulting in a high-symmetric structure with a close to tetrahedral coordination at each metal atom. In the LiCl·D2*·LiCl case the Li⁺ cations are forced out of this axis to accommodate the isopropyl groups at both ends. Another possibility is that the dimer of $[2 \cdot \text{THF}]_2$ can dissociate into monomeric units M1·LiCl. These monomers could recombine via a four-membered [MgCl]₂ ring (LiCl·D2· LiCl) or by a smaller $[LiCl]_2$ ring $(M1 \cdot (LiCl)_2 \cdot M1)$ in the center. It is also possible that LiCl dissociates as a well-known dimer³⁷ [(THF)₂Li(μ Cl)₂Li(THF)₂] Li1 to produce LiCl-free species M1, D1, or D2 or that lithium dissociates as a solvent separated ion pair $[Li(THF)_4]^+$ Li2 that would produce an SSIP ate-complex like M0, where two chlorides are coordinated to the magnesium ion (Chart 2, category B). The ⁷Li spectra at RT show a singlet at 0.18 ppm (25 °C) that shifts to lower field with decreasing temperature (0.38 ppm at -100 °C, Figure 3).



Figure 3. Superposition of ⁷Li NMR spectra³⁹ of crystalline $[2 \cdot \text{THF}]_2$ (0.1 M) redissolved in THF-*d*₈ at various temperatures. On the top: ⁷Li-spectrum of neat Li1 in THF-*d*₈ at -80 °C.

The presence of LDA, where lithium coordinates directly to the diisopropyl amide, can therefore be excluded because LDA resonates in the ⁷Li NMR experiment temperature independently at about 2.0 ppm (Figure S14). Additionally, the solvent separated cation $[\text{Li}(\text{THF})_4]^+$ Li2 can be excluded to be populated at detectable concentrations too, because it is known that the solvent separated lithium ion Li2 resonances are at a negative chemical shift (-1.1 ppm in THF).³⁸

Further support is provided by the crystal structure of $[2 \cdot \text{THF}]_2$ where the lithium cation is located close to the isopropyl groups in a middle distance of 4.49 Å to the closest CH₃-protons.²¹ This relatively close distance should be detectable in a ¹H-⁷Li-HOESY experiment,^{19e} when lithium is coordinated next to the magnesium amide. In fact, the ¹H-⁷Li-HOESY spectra show at high temperatures a cross peak between the ⁷Li signal and the CH₃ signals of species **a2**, **b2**, and **c2** (Figure S9), indicating that lithium does coordinate to Hauser base **1**. At room temperature, two different species (**a2**, 2.92/1.03 ppm, category B; and **c2**, 3.42/1.33 ppm, category A) can be identified in the ¹H NMR spectrum of [**2**·THF]₂.

(Figure 4). The MW-determination of a2 gives at all temperatures a MW that fits almost perfectly the LiCl



Figure 4. Superposition of ¹H NMR spectra of crystalline $[2 \cdot \text{THF}]_2$ (0.1 M_t – CH₃ region) redissolved in THF- d_8 at various temperatures.

coordinated monomer M1·LiCl (MW_{calc} = 419 g/mol, MW_{det} = 425 g/mol, $MW_{err} = -2\%$). As compared to the salt-free monomer M1 (a1), the metal chloride cocoordinated monomers $M1(\mu Cl)_3MgCl_2$ (f1) and $M1 \cdot LiCl$ (a2) are shifted to lower field ($\Delta \delta \approx 0.02$ ppm), due to an additional electronwithdrawing metal chloride coordinated to the aggregate. The same is true for signal c2 that appears at the dimer region of D1 (Chart 2, category A) but also low field shifted by 0.02 ppm (as compared to LiCl-free dimer D1), indicating that c2 is due to the LiCl cocoordinated dimer LiCl·D1·LiCl. At 25 °C, the determined MW of c2 (MW_{calc} = 693 g/mol, MW_{det} = 525 g/ mol, $MW_{err} = 24\%$) is smaller than expected. Yet with decreasing temperature, the MW grows significantly $(MW_{det} =$ 598 g/mol at -15 °C, 618 g/mol at -40 °C, and 635 g/mol at -50 °C) until it stops growing further at -60 °C (MW_{det} = 661 g/mol). The same trend was already observed for species MgCl₂·D1(D2)·MgCl₂ (e1) of Hauser base 1. In contrast to all other species of 2, the α -CH signal of c2 shifts like e1 to lower field with lower temperature (Figure S6). This is why a successive coordination of LiCl to D1 could again explain this behavior (25 to -40 °C: av MW_{det} = 582 g/mol, MW_{calc}(D1· LiCl) = 579 g/mol, MW_{err} = -1%).⁴⁰ Finally, between -60 and -70 °C, it is possible to determine the MW of tetranuclear dimer LiCl·D1·LiCl (c2) as two lithium chlorides coordinated to dimer D1 (MW_{calc} = 693 g/mol, MW_{det} = 661 g/mol, MW_{err} = 5%). A monomer-dimer equilibrium of a2 and c2 was already suggested by García-Álvarez and Mulvey et al.¹⁴ Concentration experiments showed that species a2 dominates at lower concentrations, whereas c2 is mostly populated at higher concentrations.¹⁴ Below -70 °C, the solubility limit is reached and signal c2 disappears. With decreasing temperature, a third species next to a2 forms (b2: 3.13/1.07 ppm, category B). Further cooling results in a shift of the oligomer equilibrium. The integral of b2 increases significantly at the expense of a2 and c2. The MW-determination of b2 offers at all temperatures a MW that fits the LiCl free dimer D2 ($MW_{calc} =$ 464 g/mol, MW_{det} = 461 g/mol, MW_{err} = 1%). The ${}^{1}\text{H}{}^{-7}\text{Li}$ HOESY experiment indicates at 0 °C the interaction of b2 with the lithium cation. However, the cross peak vanishes at lower temperatures when b2 becomes the most populated species, indicating that b2 would not coordinate strongly to LiCl. At -70 °C, next to **b**2, a fourth species **d**2 is detectable in the ¹H

experiment. At -100 °C, species d2 can be deconvoluted in the DOSY NMR experiment (Figure S31), giving a MW_{det} value that fits to LiCl coordinated species LiCl·D2·LiCl/M1·(LiCl)₂· M1 (d2) (MW_{calc} = 693 g/mol, MW_{det} = 641 g/mol, MW_{err} = 7%).⁴¹ The ⁷Li DOSY-ECC-MW-determination experiments give further information about the aggregation behavior of 2. All species produce only one signal that broadens and shifts to lower field with decreasing temperature (Figure 3). Only at 0 °C there is a small shoulder at the ⁷Li-signal, verifying that more than one Li species is present in solution of 2. From 25 to -40 °C, the ⁷Li DOSY-ECC-MW-determination gives an average, temperature-dependent MW, that reflects an average value produced by all three lithium aggregates (MW_{det} = 454-496 g/mol). Below -60 °C the MW decreases significantly without further change (MW_{det} = 382 g/mol from -60 to -100°C). This can be attributed to first LiCl·D1·LiCl (c2), precipitating from solution and the concentration of M1·LiCl (a2) decreasing significantly to very small amounts, and finally to the residual lithium chloride present in solution as the wellknown [(THF)₂Li(μ Cl)₂Li(THF)₂]-dimer Li1 (MW_{calc} = 373 g/mol, $MW_{det} = 382$ g/mol, $MW_{err} = -2\%$) that becomes the most populated Li-species at low temperatures.⁴² The formation of significant amounts of Li1 is also reflected in the chemical shift of the 7Li nucleus. With decreasing temperature the Li1 concentration increases. This is why the ⁷Li signal moves toward the chemical shift of the LiCl dimer Li1 (Figure 3).⁴³ Additionally, the signal gets very broad. The reason for that will be discussed in the following section.

DISCUSSION

At low temperature the decreased solubility could be one reason for the ⁷Li signal broadening. Additionally, we think that the formation of LiCl·D2·LiCl/M1·(LiCl)₂·M1 from the reaction of LiCl (Li1) and dimer D2 has even a bigger impact on the signal broadening, because the chemical environment of lithium in LiCl·D2·LiCl/M1·(LiCl)2·M1 differs from that in Li1. Yet, how would the reaction of Li1 with D2 occur? We already proposed the tetrameric structures LiCl·D2·LiCl and $M1 \cdot (LiCl)_2 \cdot M1$ (Chart 2) as products. In both structures, two chlorides at the magnesium atoms show a coordination number of three. Searching the Cambridge Crystallography Database⁴⁴ reveals that magnesium chlorides, with a μ -Cl₃ coordination, predominantly form cubic aggregation modes in the solid state.^{33,45} Bickelhaupt and Solà et al. have shown that the most stable isomer of methylmagnesium chloride is a T_d -symmetric (CH₃MgCl)₄ tetranuclear cluster with a cubic (MgCl)₄ core and terminal CH₃ groups at the magnesium vertexes.⁴⁶ Li halides are also known to build stable cubic tetramers [LiX]₄ (X = Cl, Br, I)⁴⁷ and even heteroleptic cubanes in the solid state (see compound 3 in Figure 5).⁴⁸ This is why the interaction of the LiCl-dimer Li1 with dimer D2 to a cubic transition state structure is most feasible. Because the ring sizes of [LiCl]₂ and



Figure 5. An equimolar reaction of $[2,6-\text{Pmp}_2\text{C}_6\text{H}_3\text{Li}]_2$ (Pmp = 2,3,4,5,6-Me₅C₆) with CdI₂ and crystallization from a saturated hexane solution at -30 °C yields crystals of cubane 3.^{48a}

[MgCl]₂ differ, a fast dissociation of the cubane back to the corresponding dimers would be plausible. Once the cubane is formed, it could open on one side to produce the already mentioned aggregates LiCl·D2·LiCl with a [MgCl]₂ core and $M1 \cdot (LiCl)_2 \cdot M1$ with a $[LiCl]_2$ core (Scheme 4). There is evidence for an opened [LiCl]4 cube already.47d To find out which of those structures is the most stable in solution, we carried out further calculations (vide infra). Additional support for the interaction of Li1 with D2 is provided by the absence of LiCl-free species M1 ($MW_{calc} = 304 \text{ g/mol}$) and homoleptic diamidomagnesium M2 ($MW_{calc} = 369$ g/mol, see Chart 1, MW_{det} (all species) = 425–661 g/mol, $MW_{err} \ge 28\%$, and $MW_{err} \ge 13\%$). At low temperatures M2 is the main species in the solution of Hauser base 1. Ashby et al. proposed for the formation of M2 a dimer-based mechanism (see Scheme 3A).^{10a} A rearrangement of dimer **aII** to an asymmetric dimer

Scheme 3. (A) Mechanism on the Formation of R_2Mg and $MgCl_2$ via an Asymmetric Dimer aIII Proposed by Ashby; and (B) In the Presence of LiCl in Various Aggregation Modes of 2, bI, bII, and bIII are Always the Best Leaving Groups, Suppressing the Formation of Homoleptic R_2Mg and $MgCl_2$ aIV^a



^aAny solvation was omitted for clarity.

aIII forms two excellent leaving groups, which dissociate into the corresponding diorgano magnesium R_2Mg and $MgCl_2$ (aIV in Scheme 3A). In the presence of LiCl, it is possible to rationalize several structures like **bI-bVIII**. Interestingly, in those structures, the LiCl coordinated monomer **bI**, the LiCldimer **bII**, and dimeric **bIII** represent always the best leaving groups (Scheme 3B), especially at low temperatures where monomeric LiCl and **M1** are highly unfavorable. Further, the cleavage of MgCl₂ or R_2Mg would be accompanied by a release of solvent separated ions that would be very unlikely because no solvent separated species are observable in solution of

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Hauser base 1 and 2.49 If the LiCl dimer Li1 (bII) would not communicate with dimer D2 (bIII), then the formation of homoleptic M2 (aIV) should be observable, which is not the case. Without LiCl, the dimer D2 is only stable above -50 °C. Yet with LiCl, it is the main species below -50 °C in Hauser base 2 most probably because of its interaction with the LiCl dimer Li1. This contrasting behavior could be reflected in the different aggregation modes of LiCl and MgCl, in THF solution. The former is predominantly dimeric, enabling a sufficient overlap of the [LiCl]₂ ring with dimer D2, while MgCl₂ is predominantly monomeric.^{30b} This is why MgCl₂ stabilizes more efficiently in a terminal coordination mode like in $M1(\mu Cl)_3MgCl_2$ and $MgCl_2 \cdot D1(D2) \cdot MgCl_2$. To address some of the open questions regarding the interaction of Li1 and D2 in THF solution, we carried out further calculations on the complexes. We started by computing the relative free energies for M1·(LiCl)₂·M1 and LiCl·D2·LiCl (depicted in Scheme 4). Our results show that the ladder and boat

Scheme 4. Proposed Interaction of the LiCl Dimer Li 1 with the Hauser Base Dimer D2



configurations, both for the M1·(LiCl)₂·M1 and for the LiCl· D2·LiCl isomers, are very close in energy (Table S31) with differences below 3 kJ/mol (at 25 °C). This is definitely within the uncertainty of the method and would indicate no particular conformational preference. However, the structures M1· (LiCl)₂·M1 with an inner [LiCl]₂ four-membered ring are favored by 42.2 kJ/mol in free energy. The next question is whether the formation of M1·(LiCl)₂·M1 is favored in solution as well. This complex can be formed, for example, by the reaction of D2 with Li1 (eq 3) or by dimerization of two M1· LiCl molecules (eq 4, Scheme 5).

The free energies for the mixed Li-Mg dimer formation were computed in a range from -90 to +25 °C, revealing that the dimerization process of D2 and Li1 (eq 3) to give M1• (LiCl)₂•M1 is slightly exergonic (-3.3 to -14.3 kJ/mol, from lower to higher temperature, Table S33). Species M1•(LiCl)₂• M1 could, as such, be formed in the solution of *Turbo*-Hauser base 2. In contrast, the dimerization of two M1•LiCl molecules (eq 4) is unfavored due to an endergonic reaction (17.3 to 3.9 kJ/mol, Table S35). Obviously such values have to be considered with certain reservations because the molarity is changing (formally, two THF molecules are released into Scheme 5. Free Energies Computed in a Range from -90 to +25 °C

D2 + Li 1 🛁	M1·(LiCl)₂·M1 + 2 THF (-3.3 to -14.3 kJ/mol)	(3)
2 M1·LiCl	M1·(LiCl) ₂ ·M1 + 2 THF (17.3 to 3.9 kJ/mol)	(4)
D2 + Li 1 🛁	LiCl·D1·LiCl + 2 THF (-58.5 to -68.1 kJ/mol)	(5)
2 M1·LiCl 🛶	LiCI·D1·LiCI + 2 THF (-37.9 to -49.9 kJ/mol)	(6)

solution during the process, so we are dealing with the formation of 3 molecules from 2 starting products). The quantum chemical result is, therefore, strongly influenced by the translation entropy. In the gas phase, the latter can be computed from the harmonic vibrational partition function. This does not apply to the solution, because solvated molecules are not free to move and possess lower translational entropy than in the gas. Our results already include a correction as suggested by Ardura et al. on the basis of a cell model for the change in translational degrees of freedom.⁵⁰ However, this little energy gain for the formation of M1•(LiCl)2•M1 stays in good agreement with our NMR experiments, showing that D2 and Li1 interact and communicate with each other but still the equilibrium is highly on the side of free D2 and Li1. In tune with the crystal structure of 2, the highest energy gain was identified for the formation of LiCl·D1·LiCl (c2) (eqs 5 and 6 in Scheme 5). Again, the reaction of D2 and Li1 is preferred over the dimerization of two M1·LiCl molecules by 19.4 kJ/ mol.⁵¹ Interestingly, the solubility limit of LiCl·D1·LiCl is reached at approximately -70 °C, providing small amounts of M1·(LiCl)₂·M1 that is thermodynamically less stable. One additional open question is the existence of cubic intermediates through which the larger mixed Li-Mg aggregates could be formed. Because there are no experimental data available on such transient species, one can only postulate about different coordination motifs. We considered four possible cubane structures Li1·D2-cube (A–D) with two available coordination sites at each Li and an additional two at each Mg site (Figure 6). Our calculations clearly identify complexes B and C as the most unstable, given the low coordination number of the Mg sites (Table S30).

A direct comparison can be in fact established between Li1-D2-cube (A) and Li1·D2-cube (B), with A being favored over B in the considered temperature range by about 10 kJ/mol. For the preferred cubane structures A and D, one might consider the associated equilibria given by the general equations:

$$D2 + Li1 \rightleftharpoons LiCl \cdot D2 - cube + nTHF$$
(7)

$$2\mathbf{M1} \cdot \mathbf{LiCl} \rightleftharpoons \mathbf{LiCl} \cdot \mathbf{D2} \cdot \mathbf{cube} + n \mathrm{THF}$$
(8)

with n = 0 for **D** and 2 for **A**. As mentioned before, the change in molarity is an obstacle in computing the free energies. By considering two equilibria simultaneously, where in one case the molarity change entropically promotes the products (**A**) and in the other the reactants (**D**), the computed free energies for dimerization according to equilibrium 7 are given in Figure 7 (Table S32). The dimerization of two **M1·LiCl** molecules to



Li 1 D2-cube (C)

Figure 6. Overview of possible cube structures of Turbo-Hauser base 2.

form Li1·D2-cube (A–D) (eq 8) is on average 19.4 kJ/mol higher in energy (Table S34).



Figure 7. Computed free energies for dimerization of D2 + Li1 to give LiCl·D2-cube (A) and LiCl·D2-cube (D).

As expected, the formation of complex D would be most favorable at lower temperatures. The crossing point is at about -55 °C. Although the equilibrium is always shifted toward the reactants, the differences can be relatively small, as low as 10.7 kJ/mol, such that the cubic intermediate should be accessible, even if not stable relative to the metal dimer complexes D2 and Li1. Our results show that the formation of small amounts of M1·(LiCl)₂·M1 (Scheme 6) is thermodynamically possible. Most probably, this complex is formed by dimerization of D2 with Li1 via a cubic intermediate, presumably Li1·D2-cube(D) at temperatures below -55 °C. Our NMR investigations show also that ⁱPr₂NMgCl·LiCl 2 does not produce SSIPs in detectable concentrations.43 Further, the proposal that the impact of LiCl on the higher reactivity of Turbo bases rests on the deaggregation of RMgX oligomers to monomers has to be revised. Most reactions of RMgX reagents proceed in THF solution, and it is known for a long time that alkyl and aryl Grignard reagents are monomeric in THF solution.⁸ Our results show that the same is true for Hauser base 1. Moreover, the reason for the lower reactivity of LiCl-free Hauser bases should be reflected in the Schlenk-equilibrium. At low temperatures, the equilibrium in THF solution is mostly shifted to the side of homoleptic diamidomagnesium R₂Mg, where the amide ligands are highly sterically hindered in



^aAt room temperature, the equilibrium is highly on the left side. At low temperature it moves significantly to the middle. The MWs were derived from ¹H and ⁷Li DOSY-ECC-MW-determinations.¹ The accuracy of this method is in the range of $MW_{err} < \pm 9\%$.³⁵

comparison with the heteroleptic RMgCl monomers and MgCl₂ cocoordinated species. These compounds that represent the most reactive species in a Hauser base solution are only present at low concentrations. This explains why it is necessary to use a large excess of Hauser bases (2-12-fold) to achieve high conversions in synthesis.^{5d} The big advantage of LiCl is the ability to shift the Schlenk-equilibrium from the homoleptic to the heteroleptic side, especially at low temperatures. The high concentration of bimetallic complexes (like monomeric M1·LiCl (a1) as well as dimeric LiCl·D1·LiCl (c2), D2 (b2), and $M1 \cdot (LiCl)_2 \cdot M1$ (d2) should provide the most influence on the reactivity, chemoselectivity, and complex-induced proximity effects (CIPE)⁵² of Turbo-Hauser bases. It is possible that this concept could also be applied to Turbo-Grignard reagents. Related studies are in progress.

CONCLUSION

Our investigation provides a deep insight into the solution structure of the Hauser base ⁱPr₂NMgCl 1 and its Turbo derivative 'Pr₂NMgCl·LiCl 2. We were able to show that their aggregation in THF- d_8 differs significantly and that the solution composition of the existing species is highly temperature dependent. Knowing the state of the equilibrium is of essential importance for every synthetic chemist, because Hauser bases as well as their Turbo analogues are used in synthesis at variable temperatures $(-75 \text{ to } 25 \degree \text{C})$. ^{5d,53} The solution structure of 1 is best represented by the common Schlenk-equilibrium (eq 1) ${}^{i}Pr_{2}NMgCl$ (M1) \Leftrightarrow $({}^{i}Pr_{2}N)_{2}Mg$ (M2) + MgCl₂, with heteroleptic M1 to be the main species at high temperatures and homoleptic M2 at low temperatures (Scheme 7A). However, our work demonstrates that in Hauser base 1, dimeric D1 and D2 are also present in the THF solution, although alkyl magnesium chlorides do not dimerize in that solvent. Therefore, the Hauser base 1 Schlenk-equilibrium has to be extended to dimeric amido bridged species D1 as well as to MgCl₂ cocoordinated species $M1(\mu Cl)_3MgCl_2$ and $[^{i}Pr_{2}NMgCl]_{x} \cdot [MgCl_{2}]_{y} \cdot [THF]_{z}$, which exist only at low temperatures, where an excess of MgCl₂ is present. Our investigations also show that the addition of LiCl to 1 has an enormous impact on the Schlenk-equilibrium. The main Scheme 7. Main Species of Hauser Base ⁱPr₂NMgCl 1 and *Turbo*-Hauser base ⁱPr₂NMgCl·LiCl 2 in THF-d₈ Solution



advantage of LiCl is the ability to shift the Schlenk-equilibrium from the homoleptic to the heteroleptic side. At room temperature, monomeric **M1·LiCl** and dimeric **LiCl·D1·LiCl** are the main species in solution of **2** (Scheme 7B), while the latter is the most populated species at high concentration (0.5– 0.6 M). Lowering the temperature below -50 °C results in the formation of an equilibrium of **D2** and the LiCl-dimer Li1 that stabilizes the heteroleptic dimer **D2** and inhibits the formation of homoleptic (ⁱPr₂N)₂Mg (M2) and MgCl₂.

EXPERIMENTAL SECTION

X-ray Analysis. Single crystals were mounted due to their sensitivity and reactivity in inert perfluorinated oil.²⁰ The X-ray data set of $[1 \cdot \text{THF}]_2$ was collected at 100(2) K on a Bruker Smart Apex II Quazar diffractometer with an INCOATEC micro source⁵⁴ equipped with mirror-monochromated Mo K_a radiation ($\lambda = 0.71073$ Å). The structure was solved by direct methods with SHELXT⁵⁵ and refined by full-matrix least-squares on F^2 for all data with SHELXL.⁵⁶ Nonhydrogen atoms were refined with anisotropic displacement parameters. All hydrogen atoms were placed in calculated positions and refined using a riding model. The structure is severely disordered and refined with distances restraints and restraints for the anisotropic displacement parameters.⁵⁷

Crystal Data for [1·THF]₂. C₂₀H₄₄Cl₂Mg₂N₂O₂, M = 464.09 g/mol, monoclinic, space group $P2_1/n$, a = 9.766(2), b = 9.648(2), c = 14.138(3) Å, $\beta = 110.00(2)^\circ$, Z = 2, V = 1251.8(4) Å³, μ (Mo K_a) = 0.327 mm⁻¹, 33 818 reflections collected, 2666 independent reflections ($R_{\text{int}} = 3.91\%$), $\theta_{\text{max}} = 26.75^\circ$, 230 parameters refined, 686 restraints used, $R_1[I > 2\sigma(I)] = 3.00\%$, $wR_2(\text{all data}) = 7.96\%$, GOF = 1.067, largest diff. peak and hole 0.285 and -0.167 e Å⁻³. Crystallographic data for compound [1·THF]₂ have been deposited at the Cambridge Crystallographic Data Centre CCDC 1423230.

Preparation of Crystalline $[1\cdot\text{THF}]_2$. PrMgCl (62.10 mmol, 31.05 mL, 2 M in THF, 1.0 equiv) was added dropwise to bulk diisopropylamine Pr_2NH (68.31 mmol, 9.60 mL, 1.1 equiv) at room temperature. Stirring overnight and removal of the solvent in vacuo yielded an off-white powder (11.68 g, 96%, considering the loss of one molecule of THF of $[1\cdot\text{THF}]_2$). To receive crystalline product, a saturated THF solution (24 mL) of $[1\cdot\text{THF}]_2$ was prepared at room temperature. After reducing the solvent, addition of 5 mL of toluene, and storage at -6 °C, $[1\cdot\text{THF}]_2$ was isolated as colorless crystals.

Preparation of Crystalline [2·THF]₂. Following the procedure suggested by García-Álvarez and Mulvey et al.,¹⁴ we synthesized the

Turbo-Hauser base $[2 \cdot \text{THF}]_2$ by reaction of donor-base-free LDA¹⁸ (0.43 g, 4.00 mmol, 1.0 equiv) with a suspension of MgCl₂ (0.38 g, 4.00 mmol, 1.0 equiv) in THF and stirred the mixture overnight at room temperature. Removing the solvent in vacuo and recrystallization (2×) in a 1:1 mixture of THF/hexane at -45 °C afforded $[2 \cdot \text{THF}]_2$ as colorless crystals.

DOSY Measurements. Dry THF- d_8 stored over 4 Å molecular sieves under argon was used. The NMR samples were prepared by dissolving crystals of [1·THF]2 or [2·THF]2 (each 0.10 M) and the DOSY reference 1-phenylnaphthalene (PhN, 0.02 M) under argon atmosphere in THF- d_8 . The diffusion coefficients of the amide species were normalized to the fixed diffusion value of the internal reference PhN (log $D_{\text{ref,fix}}(\text{PhN}) = -8.8812$; for more information, see Supporting Information section IX).¹⁷ NMR experiments were recorded on two devices: (1) Bruker Avance 400 spectrometer equipped with a observe broadband probe with z-axis gradient coil with maximum gradient strength of 57 G/cm and (2) Bruker Ascend 400 spectrometer equipped with an inverse broadband probe with zaxis gradient coil with maximum gradient strength of 51 G/cm. All spectra were acquired using 5 mm NMR tubes, which were not spun during the measurements. All DOSY experiments were performed using a double stimulated echo sequence with bipolar gradient pulses and three spoil gradients with convection compensation (dstebpg3s).⁵⁸ The duration of the magnetic field pulse gradients was adjusted for every temperature in a range of $\delta/2 = 400-3500 \ \mu s$. The diffusion time was $\Delta = 0.1$ s. The delay for gradient recovery was 0.2 ms, and the eddy current delay was 5 ms. In each PFG NMR experiment, a series of 16 spectra on 32 K data points were collected. The pulse gradients were incremented from 2% to 98% of the maximum gradient strength in a linear ramp. After Fourier transformation and baseline correction, the diffusion dimension was processed with the Topspin 3.1 software. Diffusion coefficients, processed with a line broadening of 2 Hz, were calculated by Gaussian fits with the T1/T2 software of Topspin.

Computational Details. All structures included in this study were optimized at the B3LYP-D3/def2-SVP level of theory²⁹ (the dispersion corrections were computed with Becke-Johnson type damping).59 The electronic energies were recomputed with the def2-TZVP basis set.^{29b} Solvation effects were included through the use of the COSMO continuum model⁶⁰ both in the energy and in the optimization runs. All stationary points were confirmed to be true minima on the potential energy surface through harmonic vibrational calculations. All thermodynamic correction terms were derived from the latter. To deal with the large errors associated with low-energy vibrational modes (particularly in large complexes), we applied the quasi-rigid rotor harmonic oscillator formula proposed by Grimme, with a cutoff parameter of 100 cm⁻¹. Furthermore, corrections were included to the entropy to account for the overestimation of translational freedom in solution. The latter was included taking a cell model for the change in translational degrees of freedom.⁵⁰ All reported energy values, unless otherwise noted, correspond to Gibbs free energies. All calculations were carried out with the Orca 3.0.3 program package.62

ASSOCIATED CONTENT

S Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/jacs.6b00345.

Detailed information about the NMR measurements, the DOSY-ECC-MW-determinations, and theoretical calculations (PDF)

X-ray data for compound $[1 \cdot THF]_2$ (CIF)

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Notes

The authors declare no competing financial interest.

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REFERENCES

(1) (a) Rathman, T. L.; Schwindeman, J. A. Org. Process Res. Dev. **2014**, 18, 1192–1210. (b) Mulvey, R. E.; Robertson, S. D. Angew. Chem. **2013**, 125, 11682–11700; Angew. Chem., Int. Ed. **2013**, 52, 11470–11487.

(2) Hauser, C. R.; Walker, H. G. J. Am. Chem. Soc. 1947, 69, 295-297.

(3) Eaton, P. E.; Lee, C. H.; Xiong, Y. J. Am. Chem. Soc. 1989, 111, 8016-8018.

(4) (a) Shilai, M.; Kondo, Y.; Sakamoto, T. J. Chem. Soc. Perkin Trans. 1 2001, 442–444. (b) Kondo, Y.; Yoshida, A.; Sakamoto, T. J. Chem. Soc., Perkin Trans. 1 1996, 2331–2332.

(5) (a) Melzig, L.; Rauhut, C. B.; Knochel, P. Chem. Commun. 2009, 3536–3538. (b) Rohbogner, C. J.; Wunderlich, S. H.; Clososki, G. C.; Knochel, P. Eur. J. Org. Chem. 2009, 2009, 1781–1795. (c) Clososki, G. C.; Rohbogner, C. J.; Knochel, P. Angew. Chem. 2009, 119, 7825–7828; Angew. Chem., Int. Ed. 2007, 46, 7681–7684. (d) Krasovskiy, A.; Krasovskaya, V.; Knochel, P. Angew. Chem. 2006, 118, 3024–3027; Angew. Chem., Int. Ed. 2006, 45, 2958–2961.

(6) Allan, J. F.; Clegg, W.; Henderson, K. W.; Horsburgh, L.; Kennedy, A. R. J. Organomet. Chem. **1998**, 559, 173–179.

(7) Schlenk, W.; Schlenk, W. Ber. Dtsch. Chem. Ges. B 1929, 62, 920–924.

(8) Walker, F. W.; Ashby, E. C. J. Am. Chem. Soc. 1969, 91, 3845–3850.

(9) (a) Seven, Ö.; Bolte, M.; Lerner, H.-W. Acta Crystallogr., Sect. E: Struct. Rep. Online 2013, 69, m424–m424. (b) Nesbit, M. A.; Gray, D. L.; Girolami, G. S. Acta Crystallogr., Sect. E: Struct. Rep. Online 2012, 68, m942. (c) Ren, W.; Zi, G.; Fang, D.-C.; Walter, M. D. Chem. - Eur. J. 2011, 17, 12669–12682. (d) Avent, A. G.; Caro, C. F.; Hitchcock, P. B.; Lappert, M. F.; Li, Z.; Wei, X.-H. Dalton Trans. 2004, 1567–1577. (e) Spek, A. L.; Voorbergen, P.; Schat, G.; Blomberg, C.; Bickelhaupt, F. J. Organomet. Chem. 1974, 77, 147–151. (f) Guggenberger, L. J.; Rundle, R. E. J. Am. Chem. Soc. 1968, 90, 5375–5378.

(10) (a) Ashby, E. C.; Smith, M. B. J. Am. Chem. Soc. **1964**, 86, 4363–4370. (b) Ashby, E. C.; Becker, W. E. J. Am. Chem. Soc. **1963**, 85, 118–119.

(11) Smith, M. B.; Becker, W. E. *Tetrahedron* **1967**, *23*, 4215–4227.

(12) (a) Ashby, E. C.; Nackashi, J.; Parris, G. E. J. Am. Chem. Soc.
1975, 97, 3162–3171. (b) Ashby, E. C.; Parris, G. E. J. Am. Chem. Soc.
1971, 93, 1206–1213. (c) Ashby, E. C.; Yu, S. J. Organomet. Chem.
1971, 29, 339–348.

(13) Seyferth, D. Organometallics 2009, 28, 1598-1605.

(14) Armstrong, D. R.; García-Álvarez, P.; Kennedy, A. R.; Mulvey, R. E.; Parkinson, J. A. Angew. Chem. **2010**, 122, 3253–3256; Angew. Chem., Int. Ed. **2010**, 49, 3185–3188.

(15) Because of signal overlap problems, the authors had to use internal references with inappropriate geometries. This setup produced a molecular weight (MW) determination with an error bar of approximately $\pm 30\%$.

(16) Li, D.; Keresztes, I.; Hopson, R.; Williard, P. G. Acc. Chem. Res. **2009**, 42, 270–280.

(17) Neufeld, R.; Stalke, D. Chem. Sci. 2015, 6, 3354-3364.

(18) Neufeld, R.; John, M.; Stalke, D. Angew. Chem. 2015, 127, 7100-7104; Angew. Chem., Int. Ed. 2015, 54, 6994-6998.

(19) (a) Niklas, T.; Stalke, D.; John, M. Chem. Commun. 2015, 51, 1275–1277. (b) Pöppler, A.-C.; Granitzka, M.; Herbst-Irmer, R.;

Chen, Y.-S.; Iversen, B. B.; John, M.; Mata, R. A.; Stalke, D. Angew. Chem. 2014, 126, 13498–13503; Angew. Chem., Int. Ed. 2014, 53, 13282–13287. (c) Pöppler, A.-C.; Frischkorn, S.; Stalke, D.; John, M. ChemPhysChem 2013, 14, 3103–3107. (d) Pöppler, A.-C.; Meinholz, M. M.; Faßhuber, H.; Lange, A.; John, M.; Stalke, D. Organometallics 2012, 31, 42–45. (e) Granitzka, M.; Pöppler, A.-C.; Schwarze, E. K.; Stern, D.; Schulz, T.; John, M.; Herbst-Irmer, R.; Pandey, S. K.; Stalke, D. J. Am. Chem. Soc. 2012, 134, 1344–1351. (f) Pöppler, A.-C.; Keil, H.; Stalke, D.; John, M. Angew. Chem. 2012, 124, 7963–7967; Angew. Chem., Int. Ed. 2012, 51, 7843–7846.

(20) (a) Stalke, D. Chem. Soc. Rev. **1998**, 27, 171–178. (b) Kottke, T.; Stalke, D. J. Appl. Crystallogr. **1993**, 26, 615–619.

(21) García-Álvarez, P.; Graham, D. V.; Hevia, E.; Kennedy, A. R.; Klett, J.; Mulvey, R. E.; O'Hara, C. T.; Weatherstone, S. Angew. Chem. 2008, 120, 8199–8201; Angew. Chem., Int. Ed. 2008, 47, 8079–8081.

(22) (a) Yang, K.-C.; Chang, C.-C.; Huang, J.-Y.; Lin, C.-C.; Lee, G.-H.; Wang, Y.; Chiang, M. Y. J. Organomet. Chem. 2002, 648, 176–187.
(b) Bartlett, R. A.; Olmstead, M. M.; Power, P. P. Inorg. Chem. 1994, 33, 4800–4803.

(23) Kennedy, A. R.; Mulvey, R. E.; Robertson, S. D. Dalton Trans. 2010, 39, 9091–9099.

(24) Cambridge Structural Database CSD, version 5.36 (Updated Nov 2014), chelating amido ligands have been excluded.

(25) Batsanov, A. S.; Bolton, P. D.; Copley, R. C. B.; Davidson, M. G.; Howard, J. A. K.; Lustig, C.; Price, R. D. J. Organomet. Chem. **1998**, 550, 445–448.

(26) Deuterated species diffuse like their protonated counterparts. This is why we used $MW_{calc}(THF) = 72.11$ g/mol in the calculation of all THF coordinated species. For more information, see ref 17.

(27) All MW_{det} values are dispayed for each species as an average value, derived from DOSY-ECC-MW measurements at different temperatures; see Tables S3 and S6.

(28) Rappoport, Z.; Marek, I. *The Chemistry of Organomagnesium Compounds, Part I*; John Wiley & Sons Ltd: West Sussex, England, 2008.

(29) (a) Grimme, S.; Antony, J.; Ehrlich, S.; Krieg, H. J. Chem. Phys. 2010, 132, 154104. (b) Weigend, F.; Ahlrichs, R. Phys. Chem. Chem. Phys. 2005, 7, 3297–3305. (c) Becke, A. D. J. Chem. Phys. 1993, 98, 5648–5652. (d) Becke, A. D. Phys. Rev. A: At., Mol., Opt. Phys. 1988, 38, 3098–3100. (e) Lee, C.; Yang, W.; Parr, R. G. Phys. Rev. B: Condens. Matter Mater. Phys. 1988, 37, 785–789.

(30) (a) Stern, D.; Granitzka, M.; Schulz, T.; Stalke, D. Z. *Naturforsch., B: J. Chem. Sci.* **2010**, 65b, 719–724. (b) Tammiku-Taul, J.; Burk, P.; Tuulmets, A. J. Phys. Chem. A **2004**, 108, 133–139. (31) Sakamoto, S.; Imamoto, T.; Yamaguchi, K. Org. Lett. **2001**, 3, 1793–1795.

(32) Unfortunately, below -80 °C, the signal was too low in intensity for an accurate ECC-MW-determination.

(33) Blasberg, F.; Bolte, M.; Wagner, M.; Lerner, H.-W. Organometallics 2012, 31, 1001–1005.

(34) Species $MgCl_2 \cdot D1(D2) \cdot MgCl_2$ have a highly increased molar van der Waals density (MD_w). This is why a density correction has been done on the MW_{det} value of species $MgCl_2 \cdot D1(D2) \cdot MgCl_2$; see Supporting Information section XXXIII.

(35) The deviation is calculated by $MW_{err} = [1 - MW_{det}/MW_{calc}]$. 100%, where MW_{det} is the experimentally determined and MW_{calc} is the calculated molecular weight.

(36) (a) Krasovskiy, A.; Straub, B. F.; Knochel, P. Angew. Chem.
2006, 118, 165–169; Angew. Chem., Int. Ed. 2006, 45, 159–162.
(b) Krasovskiy, A.; Knochel, P. Angew. Chem. 2004, 116, 3396–3399; Angew. Chem., Int. Ed. 2004, 43, 3333–3336.

(37) (a) Yakimansky, A. V.; Müller, A. H. E.; Van Beylen, M. *Macromolecules* **2000**, 33, 5686–5692. (b) Ming Keong, W.; Popov, A. I. *J. Inorg. Nucl. Chem.* **1972**, 34, 3615–3622. (c) Bauer, S. H.; Ino, T.; Porter, R. F. *J. Chem. Phys.* **1960**, 33, 685–691.

(38) (a) Elschenbroich, C.; Hensel, F.; Hopf, H. In Organometallchemie; Teubner, B. G., Ed.; Verlag: Wiesbaden, 2008; Vol. 6.
(b) Gornitzka, H.; Stalke, D. Angew. Chem. 1994, 106, 695–698; Angew. Chem., Int. Ed. Engl. 1994, 33, 693–695. (c) In ref 14 the authors already showed in accord with our findings that the ⁷Li chemical shift of **2** displays a positive δ value even in a wide concentration range (higher than 0.21 ppm).

(39) The 7 Li NMR chemical shifts are referenced externally to LiCl in D₂O at 0.00 ppm.

(40) This behavior could be also attributed to a fast exchange of LiCl·D1·LiCl with its lithium free counterpart D2 or to a dissociation into monomeric M1·LiCl. Both equilibria would produce a smaller effective MW.

(41) The ¹H NMR spectrum of **2** shows at -100 °C one broad signal. However, it was possible to distinguish **b2** and **d2** in the DOSY NMR experiment due to the high difference in the MW and diffusion coefficient, respectively; see Figure S31.

(42) To prove the structure of LiCl in THF, we measured anhydrous LiCl in THF- d_8 . Our ⁷Li ECC-MW-determination confirms the dimeric structure (Li1) of LiCl in THF solution (from 25 to -75° C, on av: MW_{calc} = 373 g/mol, MW_{det} = 381 g/mol, MW_{err} = -2%; see Table S7). That result stays in very good agreement with previous work from Reich et al. who categorized LiCl as a sturdy dimer (Li1) according to HMPA titrations: Reich, H. J.; Borst, J. P.; Dykstra, R. R.; Green, P. D. J. Am. Chem. Soc. 1993, 115, 8728–8741. On the other hand, we could not confirm the SSIP [Li(thf)₄]⁺Cl⁻ published by: Fernández, I.; Martínez-Viviente, E.; Breher, F.; Pregosin, P. S. Chem. - *Eur. J.* 2005, 11, 1495–1506.

(43) Additionally, a detectable formation of the SSIP $[\text{Li}(\text{thf})_4]^+$ Li2 can be excluded (MW_{calc} = 295 g/mol, MW_{det} = 382 g/mol, MW_{err} = -29%).

(44) Based on statistics from the Cambridge Structural Database CSD, version 5.36 (updated Nov. 2014). Allen, F. H. Acta Crystallogr, Sect. B: Struct. Sci. 2002, 58, 380–388.

(45) (a) Vollet, J.; Hartig, J. R.; Baranowska, K.; Schnöckel, H. Organometallics **2006**, 25, 2101–2103. (b) Casellato, U.; Ossola, F. Organometallics **1994**, 13, 4105–4108.

(46) Jiménez-Halla, J. O. C.; Bickelhaupt, F. M.; Solà, M. J. Organomet. Chem. 2011, 696, 4104-4111.

(47) (a) Brym, M.; Jones, C.; Junk, P. C.; Kloth, M. Z. Anorg. Allg. Chem. 2006, 632, 1402–1404. (b) Courtenay, S.; Wei, P.; Stephan, D. W. Can. J. Chem. 2003, 81, 1471–1476. (c) Neumann, F.; Hampel, F.; Schleyer, P. v. R. Inorg. Chem. 1995, 34, 6553–6555. (d) Hoffmann, D.; Dorigo, A.; Schleyer, P. v. R.; Reif, H.; Stalke, D.; Sheldrick, G. M.; Weiss, E.; Geissler, M. Inorg. Chem. 1995, 34, 262–269.

(48) (a) Gridley, B. M.; Blundell, T. J.; Moxey, G. J.; Lewis, W.;
Blake, A. J.; Kays, D. L. Chem. Commun. 2013, 49, 9752–9754.
(b) Sutton, A. D.; Ngyuen, T.; Fettinger, J. C.; Olmstead, M. M.;
Long, G. J.; Power, P. P. Inorg. Chem. 2007, 46, 4809–4814.
(c) Batsanov, A. S.; Davidson, M. G.; Howard, J. A. K.; Lamb, S.;
Lustig, C.; Price, D. R. Chem. Commun. 1997, 1211–1212. (d) Buttrus,
N. H.; Eaborn, C.; El-Kheli, M. N. A.; Hitchcock, P. B.; Smith, J. D.;
Sullivan, A. C.; Tavakkoli, K. J. Chem. Soc., Dalton Trans. 1988, 381–391.

(49) Grignard reagents are also known to be very weak electrolytes, see: Seyferth, D. *Organometallics* **2009**, *28*, 1598–1605.

(50) Ardura, D.; López, R.; Sordo, T. L. J. Phys. Chem. B 2005, 109, 23618-23623.

(51) Once again with a correction to the translational entropy.

(52) Whisler, M. C.; MacNeil, S.; Snieckus, V.; Beak, P. Angew. Chem. 2004, 116, 2256–2276; Angew. Chem., Int. Ed. 2004, 43, 2206–2225.

(53) Li-Yuan Bao, R.; Zhao, R.; Shi, L. Chem. Commun. 2015, 51, 6884-6900.

(54) Schulz, T.; Meindl, K.; Leusser, D.; Stern, D.; Graf, J.; Michaelsen, C.; Ruf, M.; Sheldrick, G. M.; Stalke, D. J. Appl. Crystallogr. 2009, 42, 885–891.

(55) Sheldrick, G. M. Acta Crystallogr., Sect. A: Found. Adv. 2015, 71, 3-8.

(56) Sheldrick, G. M. Acta Crystallogr., Sect. C: Struct. Chem. 2015, 71, 3–8.

(57) (a) Thorn, A.; Dittrich, B.; Sheldrick, G. M. Acta Crystallogr., Sect. A: Found. Adv. 2012, 68, 448–451. (b) Krause, L.; Herbst-Irmer, R.; Sheldrick, G. M.; Stalke, D. J. Appl. Crystallogr. 2015, 48, 3–10.

- (c) Müller, P.; Herbst-Irmer, R.; Spek, A. L.; Schneider, T. R.; Sawaya, M. R. Crystal Structure Refinement-A Crystallographer's Guide to SHELXL; Oxford University Press: Oxford, England, 2006.
- (58) (a) Jerschow, A.; Müller, N. J. Magn. Reson. 1997, 125, 372– 375. (b) Jerschow, A.; Müller, N. J. Magn. Reson., Ser. A 1996, 123, 222–225.
- (59) (a) Johnson, E. R.; Becke, A. D. J. Chem. Phys. 2006, 124,
- 174104. (b) Becke, A. D.; Johnson, E. R. J. Chem. Phys. 2005, 123, 154101. (c) Johnson, E. R.; Becke, A. D. J. Chem. Phys. 2005, 123, 024101.

(60) Klamt, A.; Schüürmann, G. J. Chem. Soc., Perkin Trans. 2 1993, 799–805.

(61) Grimme, S. Chem. - Eur. J. 2012, 18, 9955-9964.

(62) Neese, F. WIREs Comput. Mol. Sci. 2012, 2, 73-78.